



Part III Symmetry and Bonding

Chapter 3 Direct Products

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3. Direct products

- In this chapter we will learn how to find *the symmetry of a product of two or more functions*.

This is extraordinarily important!

Recall those integrals we used before:

$$S_{ij} = \int \psi_i^* \psi_j d\tau \quad \beta = \int s_a \hat{H} s_b d\tau$$

3.1 Introduction

- From the C_{2v} character table, we know that *the function x transforms like B_1* whereas *the function y transforms like B_2* . Then how does the function xy transform?

- This is already given in the table.

The function xy transforms like A_2 .

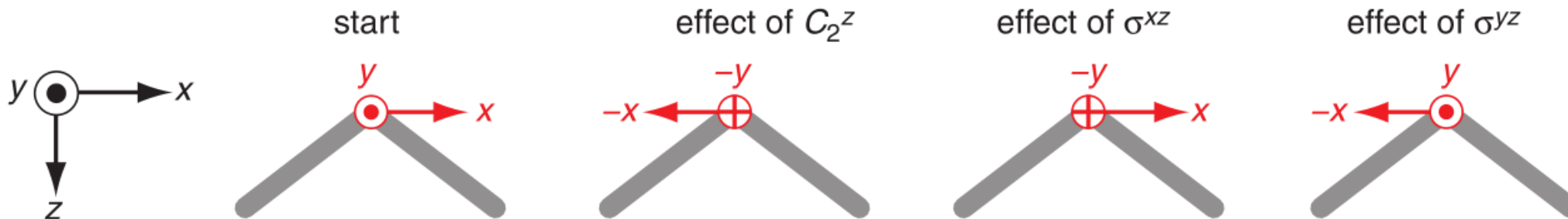
- How can we actually work this out?

C_{2v}	E	C_2^z	σ^{xz}	σ^{yz}		
A_1	1	1	1	1	z	$x^2; y^2; z^2$
A_2	1	1	-1	-1	R_z	xy
B_1	1	-1	1	-1	x R_y	xz
B_2	1	-1	-1	1	y R_x	yz



3.1 Direct products introduction

- Use the function xy as a basis to form the corresponding representation of C_{2v} , which will just be *a set of numbers*, i.e., these numbers are the *characters*.



x	$Ex = (+1)x,$	$C_2^zx = (-1)x,$	$\sigma^{xz}x = (+1)x,$	$\sigma^{yz}x = (-1)x,$	$(1,-1,1,-1) B_1$
y	$Ey = (+1)y;$	$C_2^zy = (-1)y;$	$\sigma^{xz}y = (-1)y,$	$\sigma^{yz}y = (+1)y,$	$(1,-1,-1,1) B_2$
xy	$E(xy) = (+1)xy;$	$C_2^zxy = (+1)xy;$	$\sigma^{xz}xy = (-1)xy,$	$\sigma^{yz}xy = (-1)xy,$	$(1,1,-1,-1) A_2$

C_{2v}	E	C_2^z	σ^{xz}	σ^{yz}	
A_1	1	1	1	1	$z \quad x^2; y^2; z^2$
A_2	1	1	-1	-1	$R_z \quad xy$
B_1	1	-1	1	-1	$x \quad R_y \quad xz$
B_2	1	-1	-1	1	$y \quad R_x \quad yz$

➤ xy transforms like A_2 .



- The characters for xy are simply found by multiplying together the characters for the IR B_1 , which is how x transforms, and for the IR B_2 , which is how y transforms, operation by operation:

$$\underbrace{(1, -1, 1, -1)}_{B_1 (x)} \otimes \underbrace{(1, -1, -1, 1)}_{B_2 (y)} = (1 \times 1, -1 \times -1, 1 \times -1, -1 \times 1) \equiv \underbrace{(1, 1, -1, -1)}_{B_1 \otimes B_2 = A_2}$$

This kind of multiplication is called the *direct product*: $B_1 \otimes B_2 = A_2$.

- To take another example, if we wanted to know how xz transforms:

$$\underbrace{(1, -1, 1, -1)}_{B_1 (x)} \otimes \underbrace{(1, 1, 1, 1)}_{A_1 (z)} = (1 \times 1, -1 \times 1, 1 \times 1, -1 \times 1) \equiv \underbrace{(1, -1, 1, -1)}_{B_1 \otimes A_1 = B_1}$$

Thus xz transforms like B_1 .

C_{2v}	E	C_2^z	σ^{xz}	σ^{yz}		
A_1	1	1	1	1	z	$x^2; y^2; z^2$
A_2	1	1	-1	-1	R_z	xy
B_1	1	-1	1	-1	x	R_y xz
B_2	1	-1	-1	1	y	R_x yz



3.2 Direct product of one-dimensional irreducible representations

◆ One-dimensional *IRs* are those with character *1* under *the operation E*, and always denoted by *the labels A and B*.

全对称不可约表示

◆ In any group there is always the *totally symmetric IR* with all of the characters being *+1*.

◆ For the *i*th one-dimensional *IR*, $\Gamma^{(i)}$, of a group, the following properties apply:

1) The direct product of this *IR* with the *totally symmetric IR*, $\Gamma^{tot. sym.}$, gives this *IR*,

$$\Gamma^{(i)} \otimes \Gamma^{tot. sym.} = \Gamma^{(i)}$$

2) The direct product of a *one-dimensional IR* with itself gives the *totally symmetric IR*

$$\Gamma^{(i)} \otimes \Gamma^{(i)} = \Gamma^{tot. sym.}$$

C_{2v}	E	C_2^z	σ^{xz}	σ^{yz}		
A_1	1	1	1	1	z	$x^2; y^2; z^2$
A_2	1	1	-1	-1	R_z	xy
B_1	1	-1	1	-1	x	R_y
B_2	1	-1	-1	1	y	R_x



3.3 Direct product of two-dimensional irreducible representations

- Two-dimensional **IRs** have character **2** under *the identity operation*, and are always denoted by *a label E*. (e.g., **E** IR in C_{3v})

C_{3v}	E	$2C_3^z$	$3\sigma_v$	
A_1	1	1	1	z $x^2 + y^2; z^2$
A_2	1	1	-1	R_z
E	2	-1	0	(x, y) (R_x, R_y) $(xz, yz); (x^2 - y^2, 2xy)$

- Property 1** from the previous section still applies. For example, if we take the direct product $A_1 \otimes E$ we obtain **E**.

$$E \otimes E \quad 4 \quad 1 \quad 0 \quad = E \oplus A_1 \oplus A_2$$

$$\underbrace{(1, 1, 1)}_{A_1} \otimes \underbrace{(2, -1, 0)}_E = (1 \times 2, 1 \times -1, 1 \times 0) \equiv \underbrace{(2, -1, 0)}_E \quad \underbrace{(2, -1, 0)}_E \otimes \underbrace{(2, -1, 0)}_E = (2 \times 2, -1 \times -1, 0 \times 0) \equiv (4, 1, 0)$$

- Property 2** does not apply. If we compute $E \otimes E$, we find $E \otimes E = E \oplus A_1 \oplus A_2$
- Modified version of property 2:** The direct product of an IR with itself *contains* the *totally symmetric IR*. This trend holds for higher-dimensional **IRs**.



3.4 Further points

- How does xyz transform in the group C_{2v} ? Consider the triple direct product:

$$\underbrace{B_1}_x \otimes \underbrace{B_2}_y \otimes \underbrace{A_1}_z = \underbrace{A_2}_{B_1 \otimes B_2} \otimes \underbrace{A_1}_z = A_2.$$

Thus xyz transforms as A_2 .

C_{2v}	E	C_2^z	σ^{xz}	σ^{yz}	
A_1	1	1	1	1	z $x^2; y^2; z^2$
A_2	1	1	-1	-1	R_z xy
B_1	1	-1	1	-1	x R_y xz
B_2	1	-1	-1	1	y R_x yz

- The direct product is commutative and distributive.

i.e. $B_1 \otimes B_2 = B_2 \otimes B_1$ and $(B_1 \otimes B_2) \otimes A_1 = B_1 \otimes (B_2 \otimes A_1)$.

- Simple numbers** (scalars) transform as the *totally symmetric IR*, as a number is *unaffected* by any symmetry operation.

Ex.13



3.5 Summary

- If *two functions* transform as the **IRs** $\Gamma^{(i)}$ and $\Gamma^{(j)}$, respectively, then *their product* transforms as the *direct product* of the two **IRs** $\Gamma^{(i)} \otimes \Gamma^{(j)}$.
- The *direct product* is found by *multiplying the characters* of the two **IRs** for each symmetry operation: $(a, b, c, \dots) \otimes (p, q, r, \dots) = (a \times p, b \times q, c \times r, \dots)$
- The *totally symmetric IR*, $\Gamma^{tot. sym.}$, has character **+1** for all operations.
- For any **IR** $\Gamma^{(i)}$: $\Gamma^{(i)} \otimes \Gamma^{tot. sym.} = \Gamma^{(i)}$.
- For any one-dimensional **IR**: $\Gamma^{(i)} \otimes \Gamma^{(i)} = \Gamma^{tot. sym.}$.
- For any higher-dimensional **IR** the result of the product $\Gamma^{(i)} \otimes \Gamma^{(i)}$ contains $\Gamma^{tot. sym.}$.
- *Scalars (numbers)* transform as the *totally symmetric IR*.